

# Upscale and Downscale Error Growth

Dale Durran

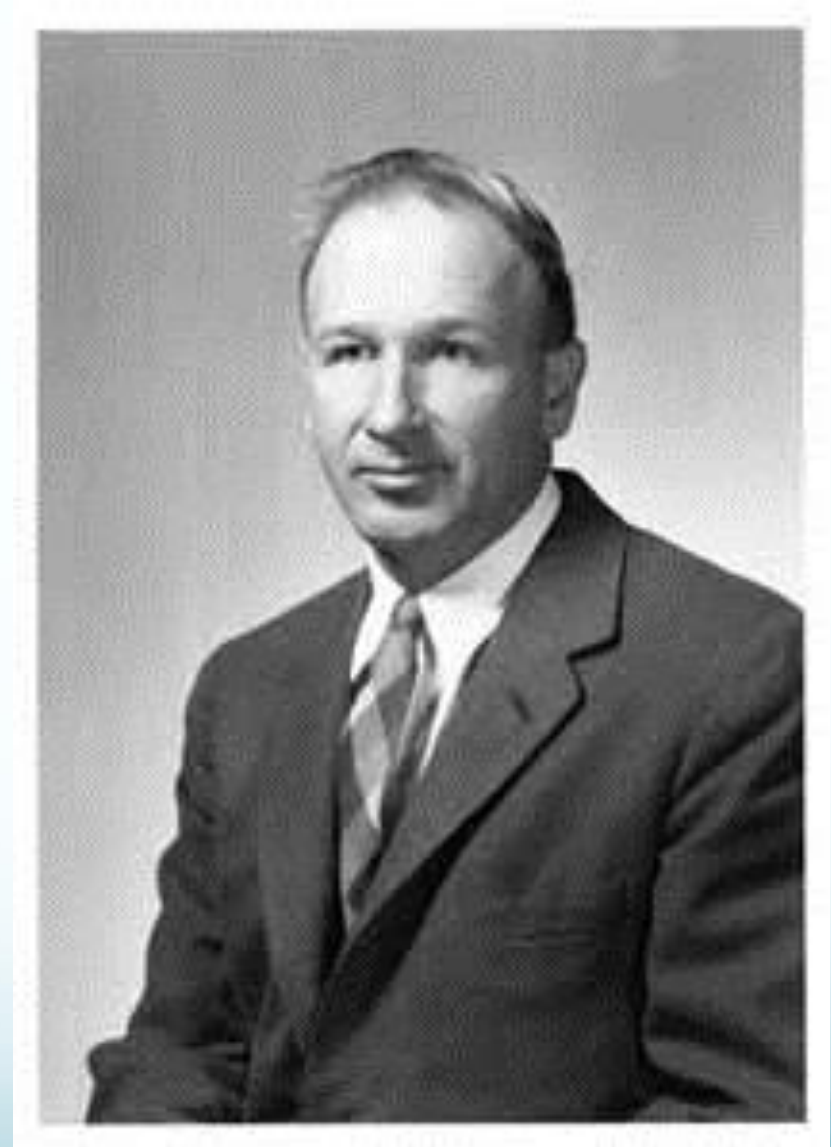
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# Edward Lorenz

## Flavors of predictability

- Chaos
  - Sensitive dependence to initial conditions, bounded phase space
- Butterfly effect
  - Finite limit on predictability due to upscale cascade of errors



# Which Lorenz equation set has a “butterfly effect”?

Nonlinear

$$\begin{aligned}\dot{x} &= \sigma(y - x) \\ \dot{y} &= x(r - z) - y \\ \dot{z} &= xy - bz\end{aligned}$$

**Chaos** (Lorenz, 1963)

Linear (with nonlinear saturation)

$$\ddot{Z}_k = \sum_{l=1}^{21} C_{k,l} Z_l$$

**Butterflies** (Lorenz, 1969)

*“The predictability of a flow which possesses many scales of motion”*

$$\ddot{Z}_k = \sum_{l=1}^n C_{k,l} Z_l$$

- $Z_k$  ensemble mean error KE at wavenumber  $k$
- $C_{k,l}$  strength of interaction of error KE at  $l$  with the background flow to force errors at  $k$
- $C_{k,l}$  and speed of upscale error propagation depend on the slope of the background KE spectrum

# Upscale error propagation: scaling argument

- $E(k)$ : background KE per unit wavenumber ( $\text{m}^3 \text{s}^{-2}$ )
- Dimensional analysis: time scale is  $\tau(k) = k^{-3/2} E^{-1/2}$ 
  - Assumed to be the time required for errors propagate from  $k$  to  $k/2$
- $T_{\text{up}}$ : time for error propagation upscale from  $2^N k_L$  to  $k_L$  is sum of time scales  $\tau(k)$  over geometrically increasing wavelengths

$$\tau(2^N k_L) + \tau(2^{N-1} k_L) + \dots + \tau(2k_L)$$

Palmer et al., 2014: The real butterfly effect

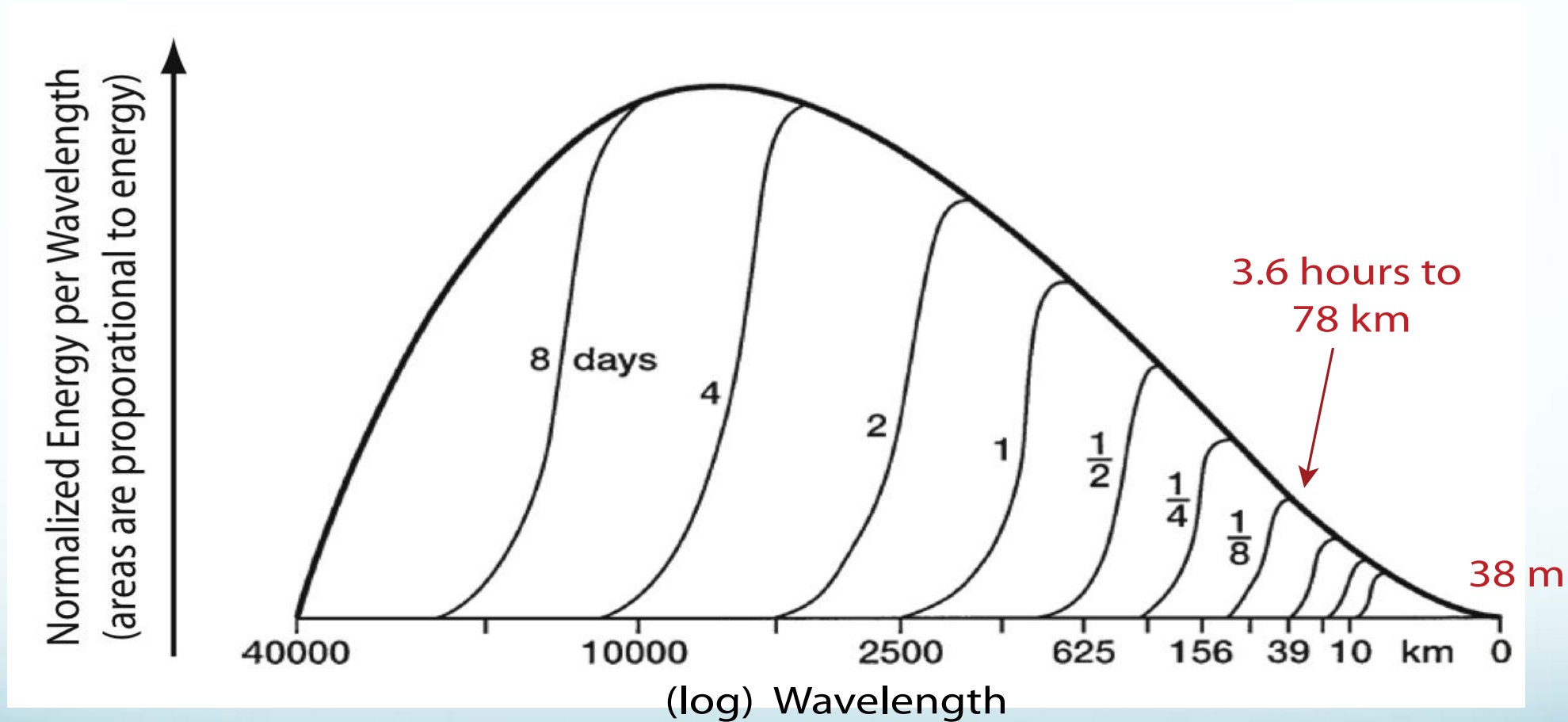
# Upscale error propagation: scaling argument

- For background KE spectrum proportional to  $k^{-p}$ ,

$$\tau(2^n k_L) \propto [2^{(p-3)/2}]^n$$

- As  $N \rightarrow \infty$ ,  $T_{up}$  converges to a **finite value** if  $p < 3$ .
- Finite limit on predictability when  $p < 3$ .

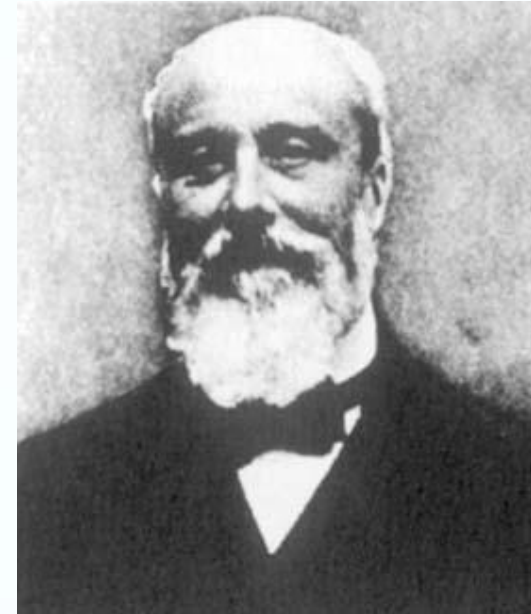
# Lorenz: 1969. Errors propagate upscale in turbulent flows with a $k^{-5/3}$ KE spectrum.



Lorenz, 1969: The predictability of a flow which possesses many scales of motion. *Tellus*, 21, 289-307.

# Butterflies before Lorenz

W.S. Franklin in 1898 review of book by Pierre Duhem



“An infinitesimal cause produces a finite effect. Long range detailed weather prediction is therefore impossible... the **flight of a grasshopper** in Montana may turn a storm aside from Philadelphia to New York!”

G. Vallis, 2006, p. 372

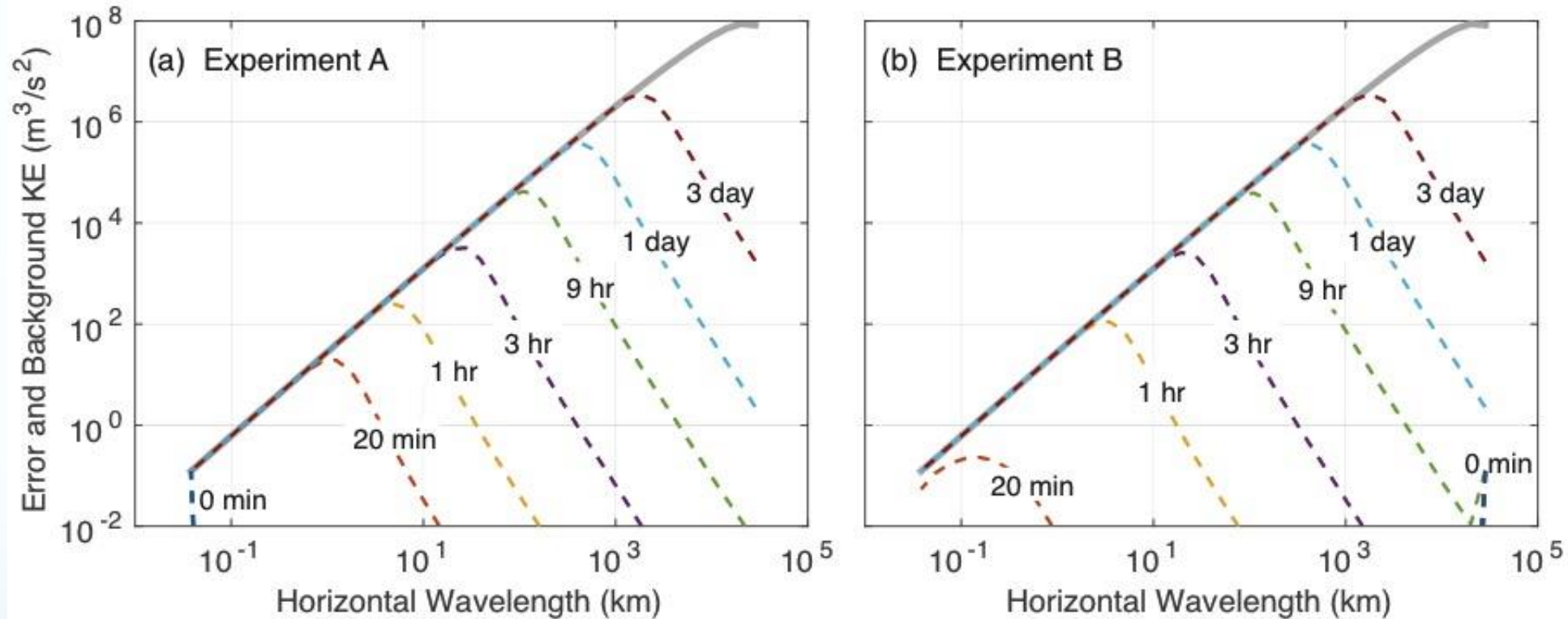


Should we associate the butterfly with small-amplitude or small-scale perturbations?

# Consider two *different* questions

- Is upscale error growth important?
  - (even if it is not exactly a “spectral cascade”)
- Given initial errors of *fixed absolute magnitude*, does their *horizontal scale* influence predictability?

# Lorenz, 1969: Experiments A & B



“Evidently when the initial error is small enough, its spectrum has little effect upon the range of predictability.”

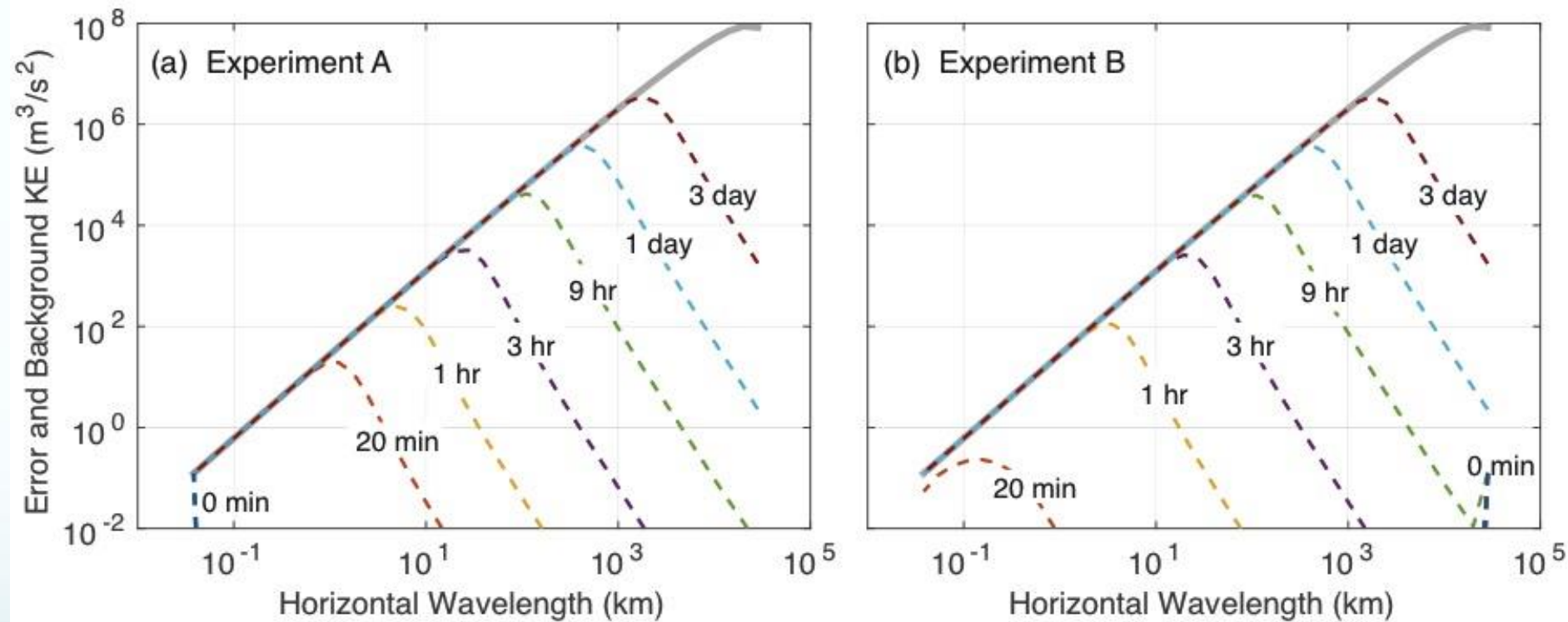
Implications of Experiment B were largely overlooked.

# Anthes, 1985: The large scale makes mesoscale predictability possible.

## *Downscale error quenching*

- Estimates of mesoscale predictability from classical turbulence theory are too pessimistic.
  - Physical forcing at the earth's surface, such as mountains, may contribute to extended predictability.
  - Mesoscale phenomena, such as fronts, can evolve from purely large-scale initial conditions.
- Fine-resolution forecast models rely on large-scale and surface forcing to create small-scale features that cannot be initialized based directly on observations.
- *How accurate must the large-scale forecast be?*

# Small *relative errors* in the large-scales can destroy predictability.



Relative error in the velocity  
100% at  $\lambda=38$  m.

Relative error in the velocity  
1% at  $\lambda=28,000$  km.

# The large scale must be known with extreme accuracy to correctly

- Forecast the *mesoscale* distribution of orographic precipitation (Nuss and Miller, 2001)
- Forecast downslope winds (Reinecke and Durran, 2009)
- Differentiate between lowland rain and snow in Pacific Northwest winter storms (Durran et al., 2013)

*The large-scale does exert a strong control, but small errors in the large scale often interfere with mesoscale predictability.*

# On the scale of initial errors

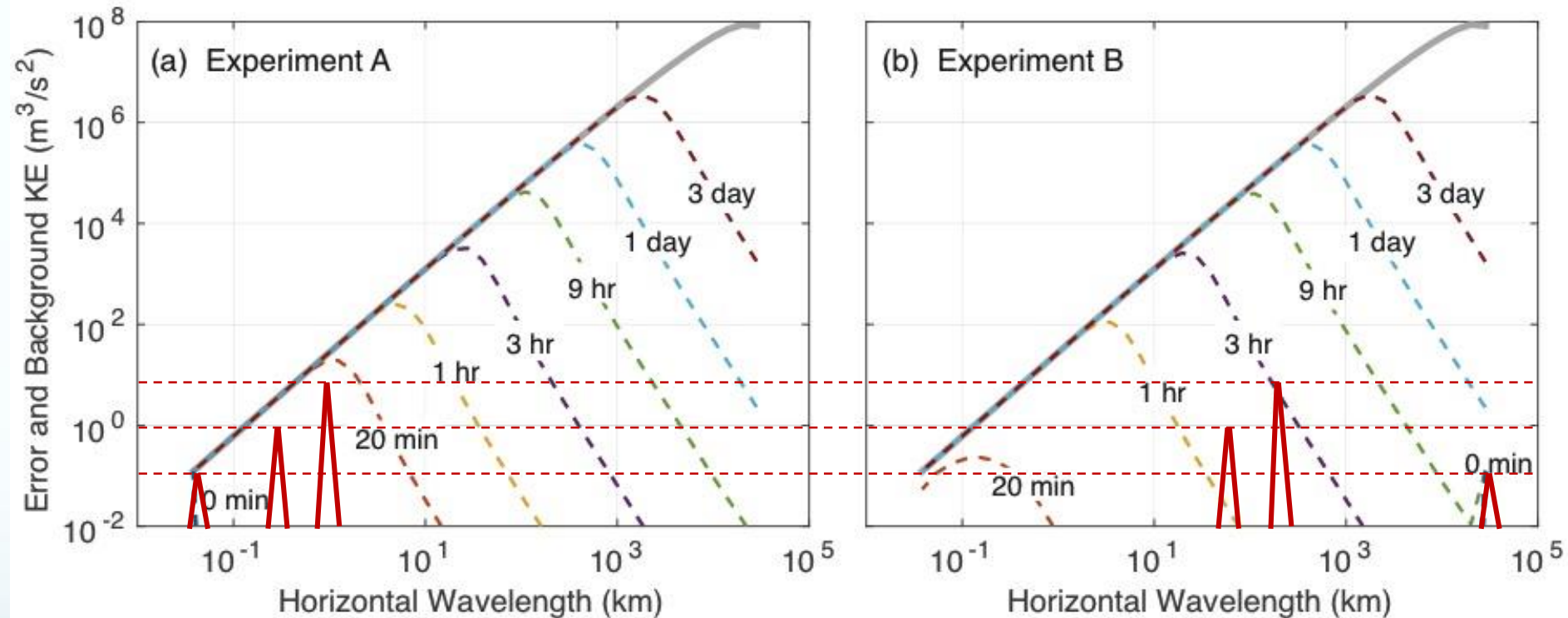
- Lorenz (1969): “We have proposed that certain formally deterministic fluid systems possessing many scales of motion may be observationally indistinguishable from indeterministic systems, in that **they possess an *intrinsic range of predictability* which cannot be lengthened by reducing the error of observation to any value greater than zero.**”
- Limits from intrinsic predictability become apparent as saturated errors appear at progressively larger scales.
- **Scale of the initial error is largely irrelevant** (for wavelengths < 400 km)

# More implied by Lorenz 1969

- Small-amplitude large-scale errors rapidly propagate downscale (for  $k^{-5/3}$  background KE spectrum) and then propagate back up scale.
- *Net effect appears as if the error originated in the smallest scales.*
- *Rate of upscale error growth does determine the theoretical limit of intrinsic predictability.*



# Initial amplitude, not initial scale

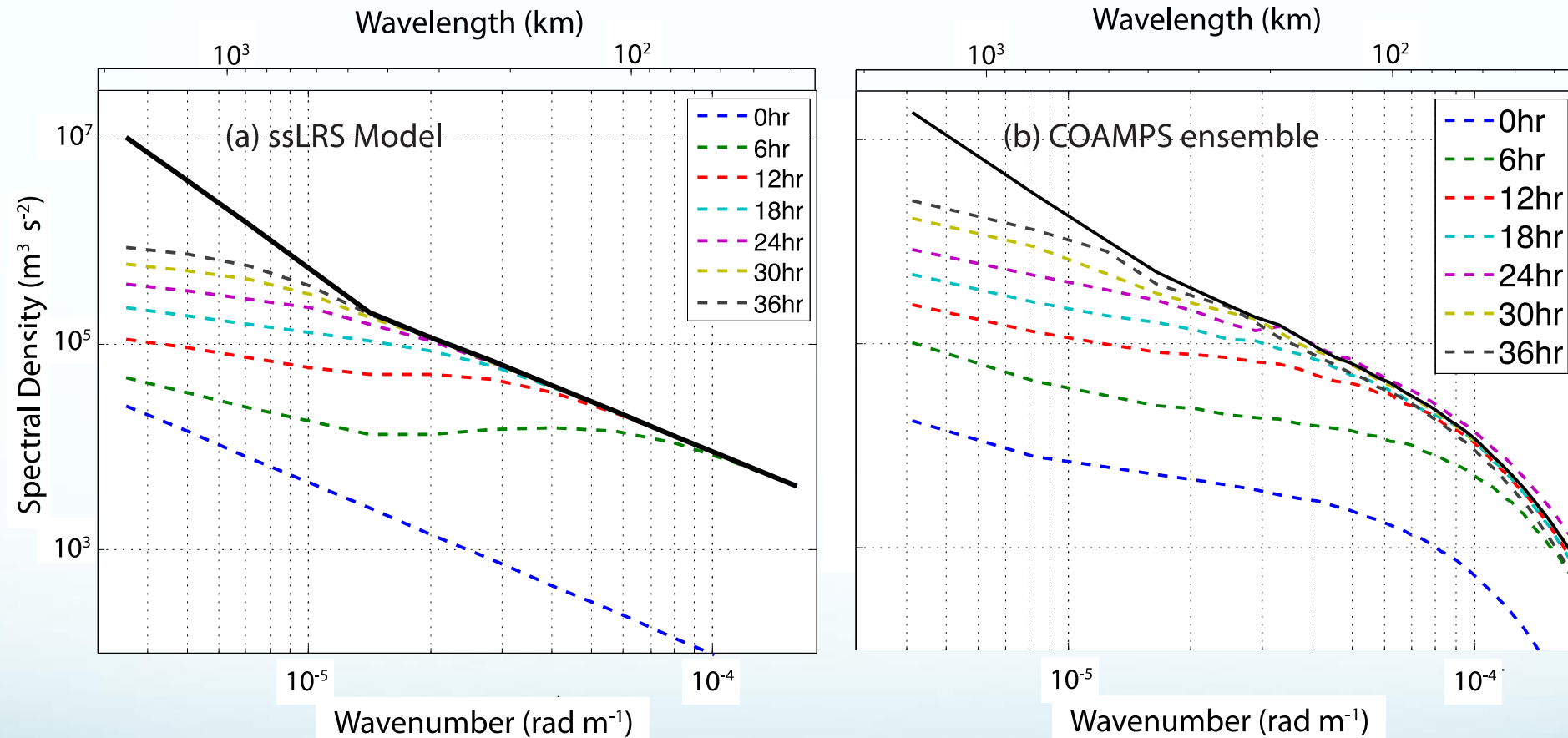


- After downscale propagation, up-scale error growth begins at a smaller scale if the error is smaller amplitude.
  - Faster eddy turnover time on smaller scales.

# How relevant is the Lorenz model?

- It does not include
  - Baroclinic instability
  - Deep convection
  - Inhomogeneity and nonstationarity
- Nonlinear effects are incorporated only crudely.
- Incorrectly assumed  $k^{-5/3}$  slope for the background KE spectrum at large-scales.
- *Nevertheless*, when given appropriate initial errors ( $\sim 1\%$  relative errors at all scales), it predicts error growth similar to that in 100-member COAMPS forecasts of east-coast snowstorms.

# Lorenz model compared with 5-km-resolution COAMPS ensemble simulations of east-coast snowstorm



*Error growth is more up-amplitude than up-scale*

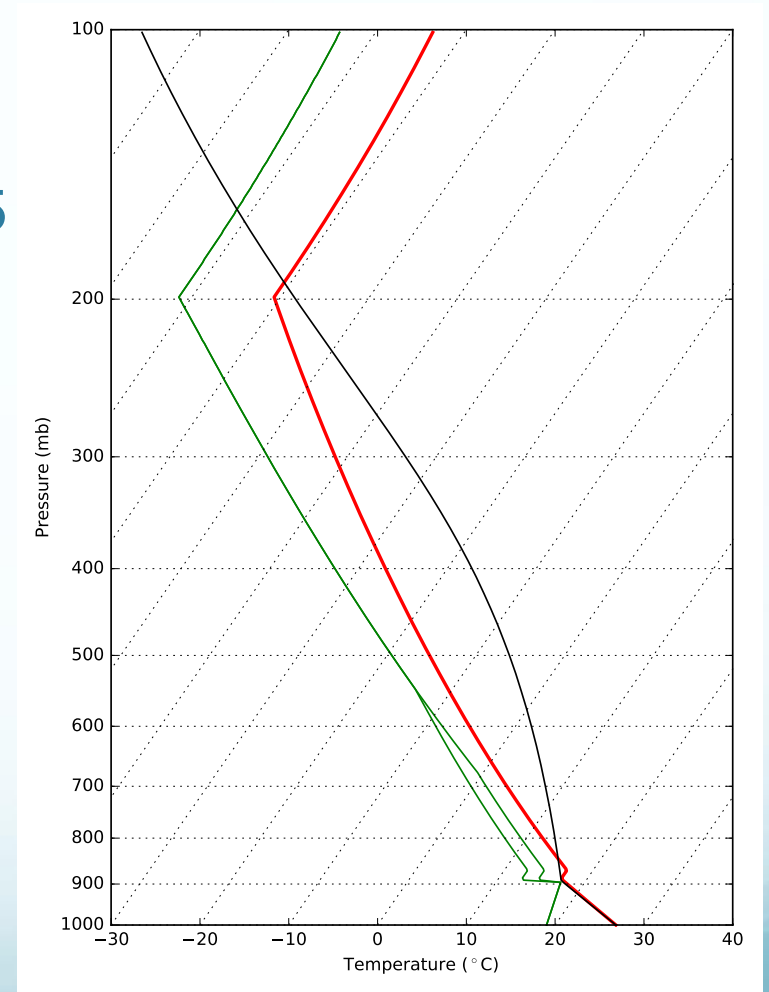
# Does the scale of initial mesoscale errors matter?

- Idealized convective systems
- Actual convective cases
- Idealized moist baroclinic instability



# 20-member ensemble simulations of deep convection

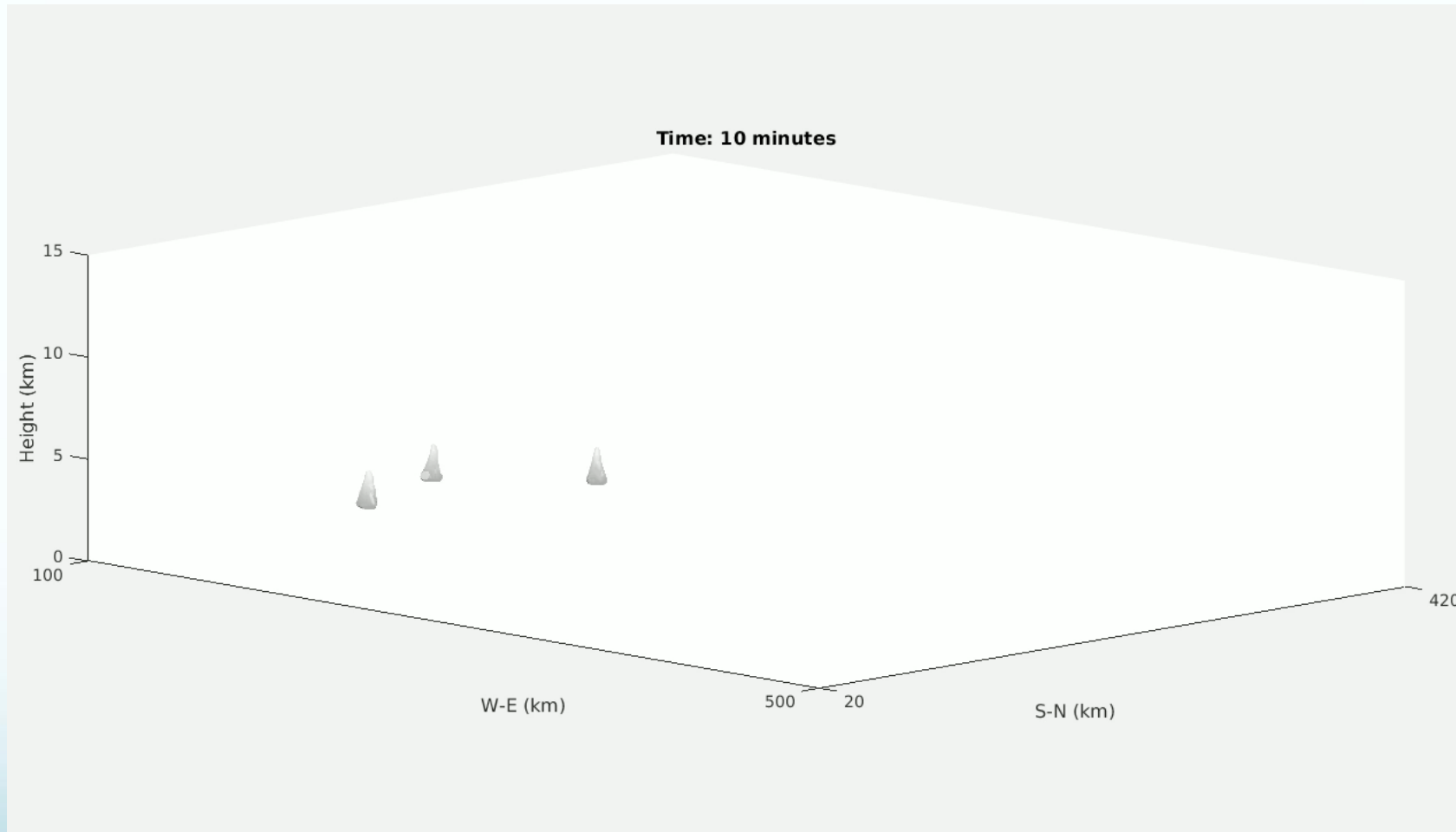
- Modified Weismann and Klemp 1983 idealized sounding
  - Unidirectional shear from 0 to 10, 20 or 30 m s<sup>-1</sup> over 5 km
  - Shear favors organization of the convection into a squall line
- 512 km x 512 km doubly periodic horizontal domain
  - Facilitates spectral analysis
  - 1 km horizontal, 40 to 500 m vertical grid spacing
  - Surface friction, but no surface heat fluxes
- Coriolis force neglected



# Ensemble strategy

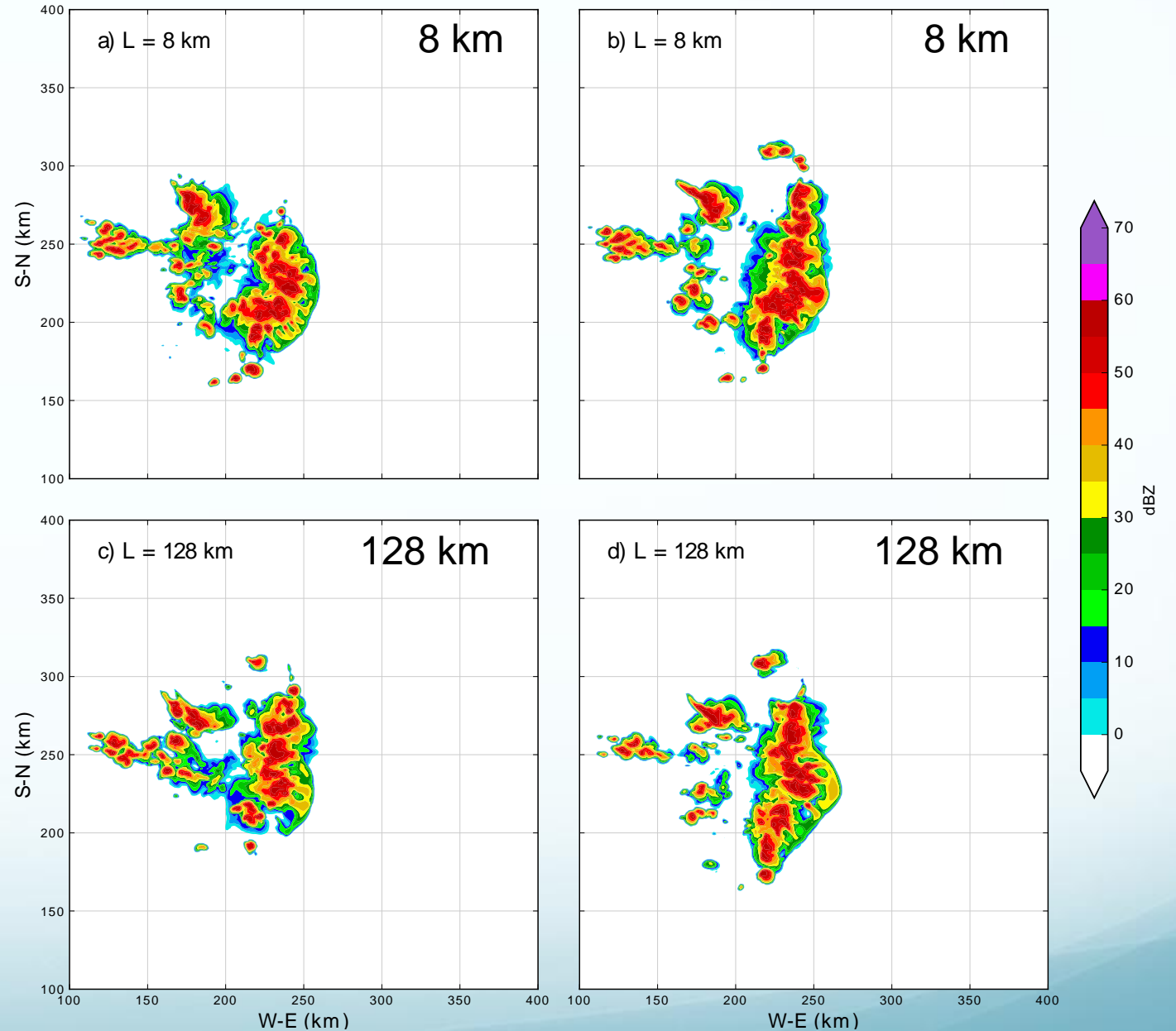
- All members initialized with 3 identical 2 K warm bubbles in the *same* location
- *Different* background perturbations among members in the near-surface *moisture* field
  - Monochromatic square wave in horizontal, random phase
    - Small-scale ensemble: *wavelength* 8 km
    - Large-scale ensemble: *wavelength* 128 km
  - Perturbation amplitude of  $0.1 \text{ g kg}^{-1}$
  - 1-km e-folding decay scale away from the surface
- Simulate for 6 hours

# Evolution of one member: $w$ and cloud



# Variability among ensemble members

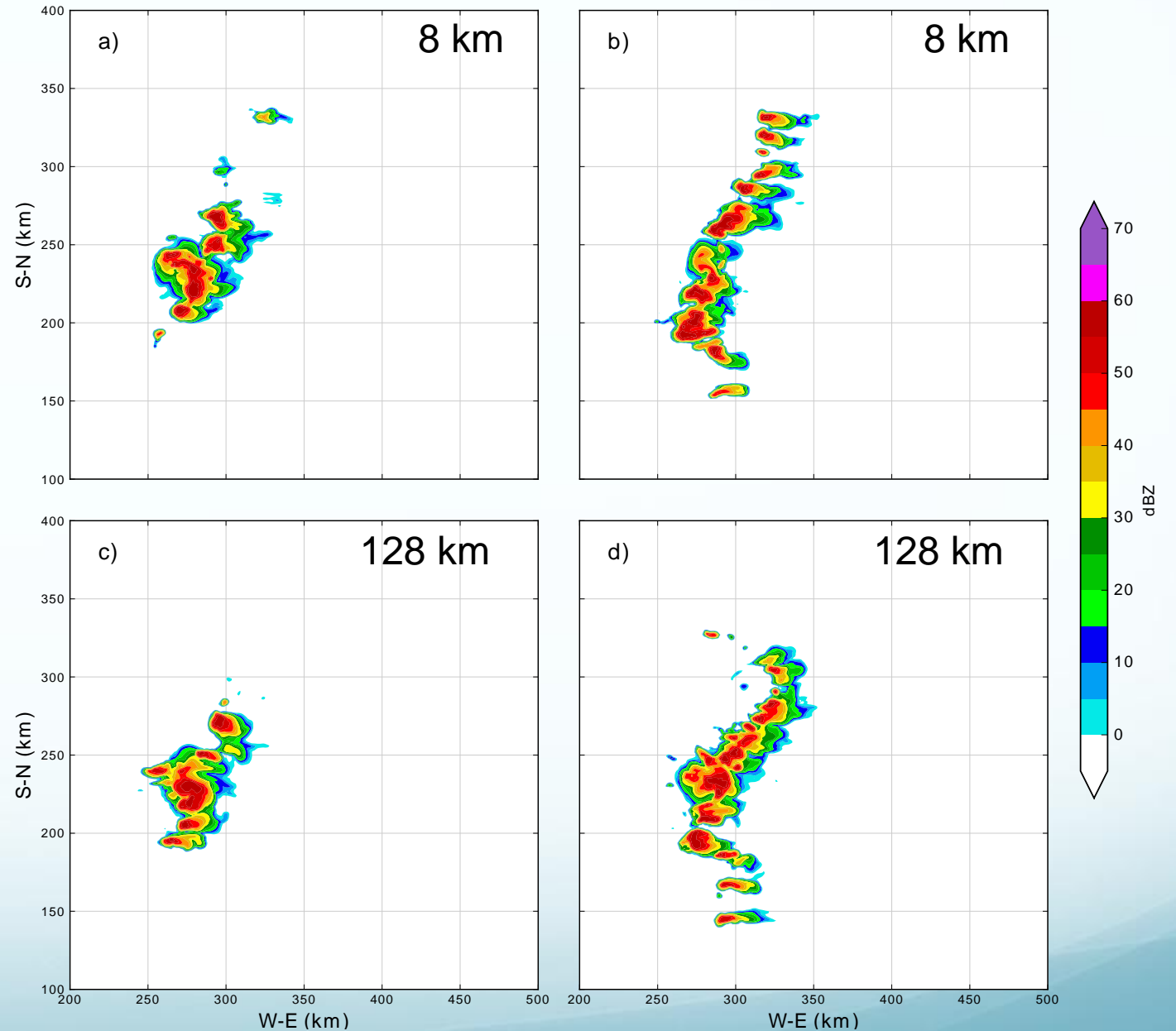
- Synthetic reflectivity (colors) and anvil-level cloud water (gray)
- T=5 hours
- 10 m s<sup>-1</sup> shear





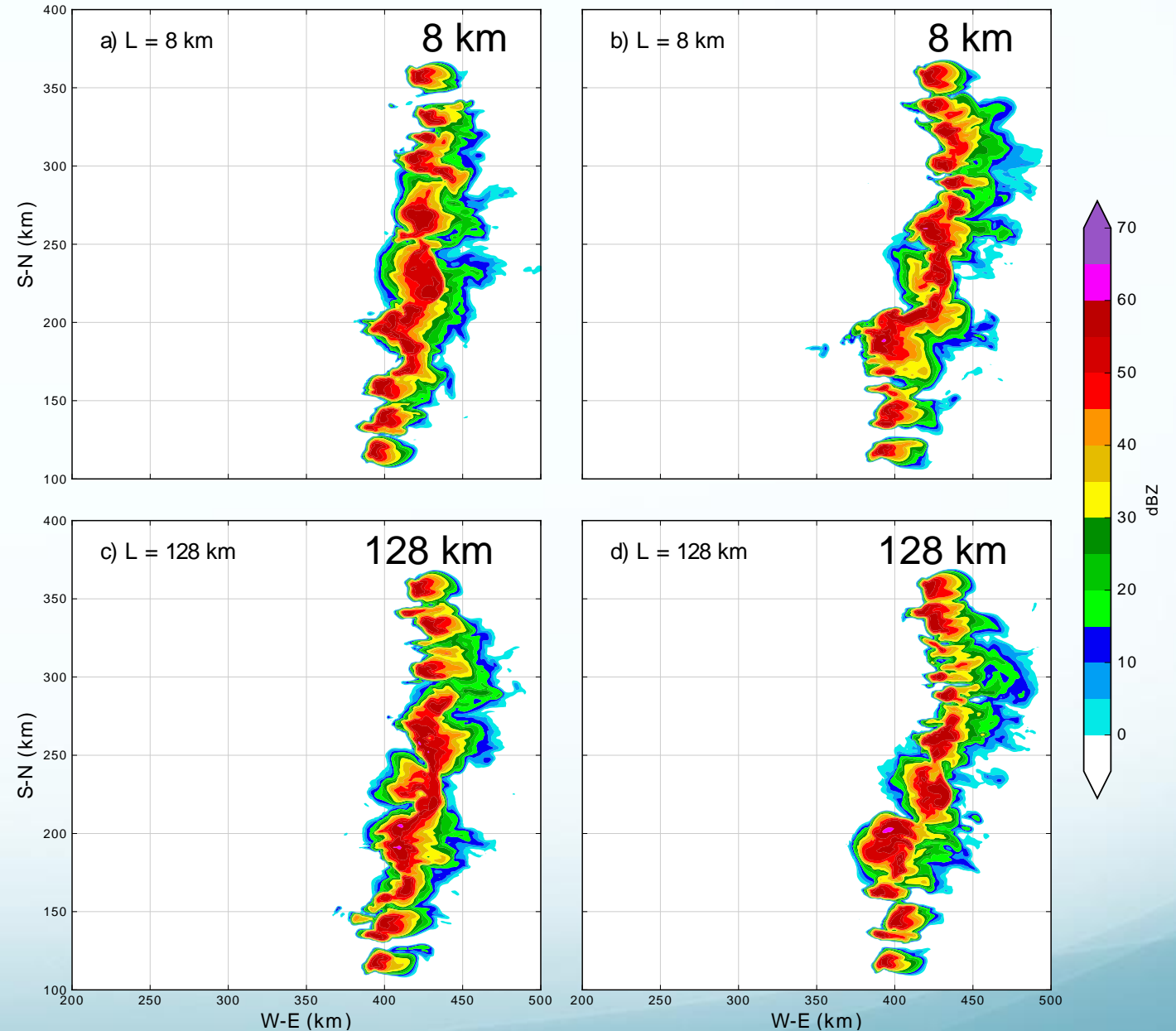
# Variability among ensemble members

- Synthetic reflectivity and anvil-level cloud water
- T=5 hours
- $20 \text{ m s}^{-1}$  shear

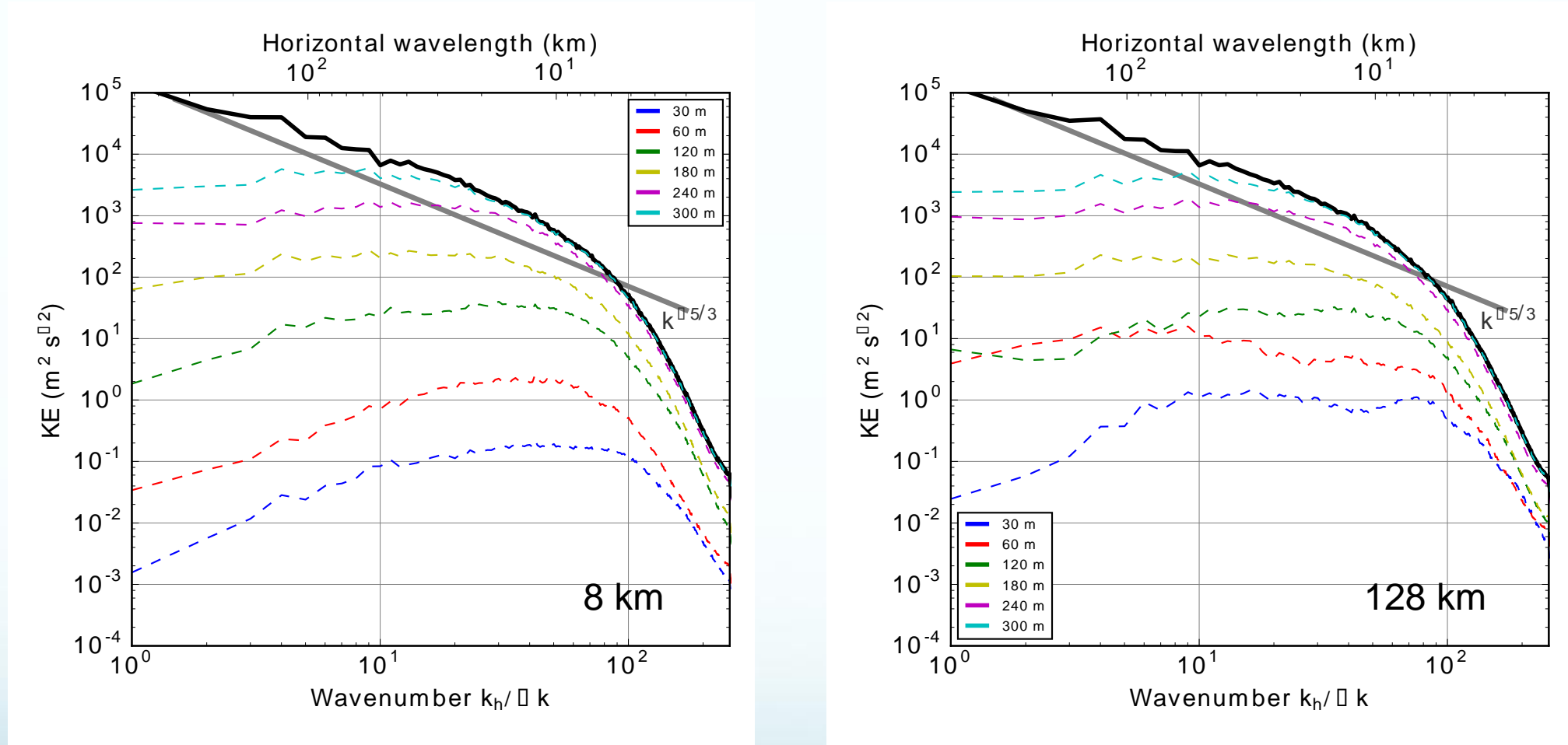


# Variability among ensemble members

- Synthetic reflectivity and anvil-level cloud water
- T=5 hours
- 30 m s<sup>-1</sup> shear

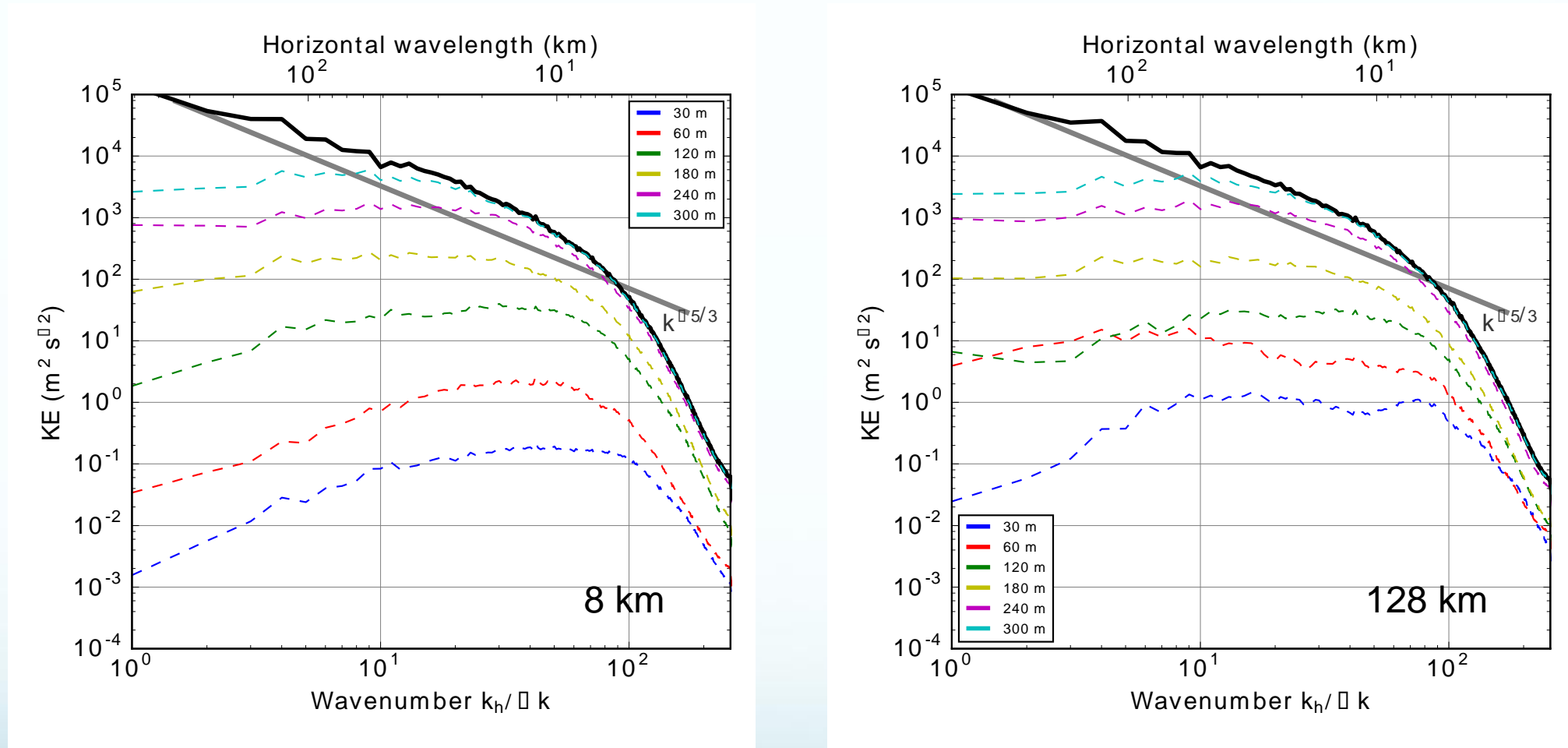


# Ensemble mean KE' (error) and KE for 20 m s<sup>-1</sup> shear



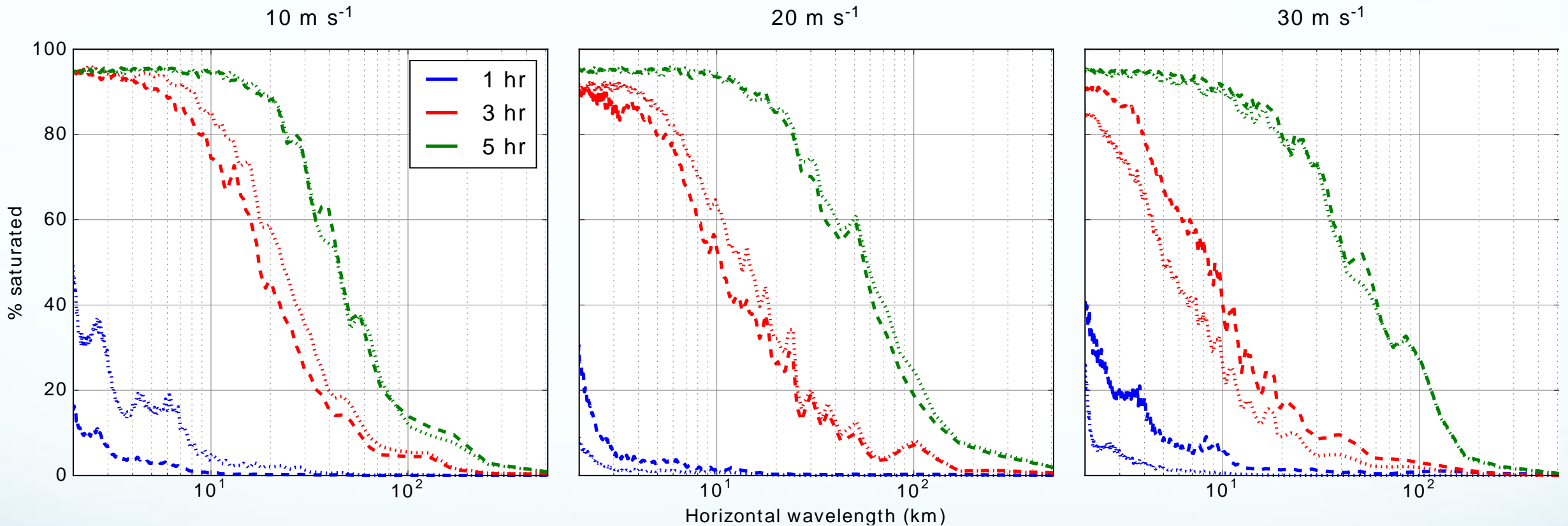
KE' dashed; black line is KE at 5 hr, gray line is observed background KE

# Ensemble mean KE' (error) and KE for 20 m s<sup>-1</sup> shear



***Error growth is up-amplitude, not an up-scale cascade!***

# Error saturation ( $KE'/KE$ ) in layer $10 < z < 12$ km



- 8-km ensemble dotted; 128 is dashed
- *By 5 hours the degree of saturation is independent of initial scale*
  - 8-km ensemble starts saturating faster for weak shear
  - 128-km ensemble starts faster for strong shear

# Intrinsic Predictability

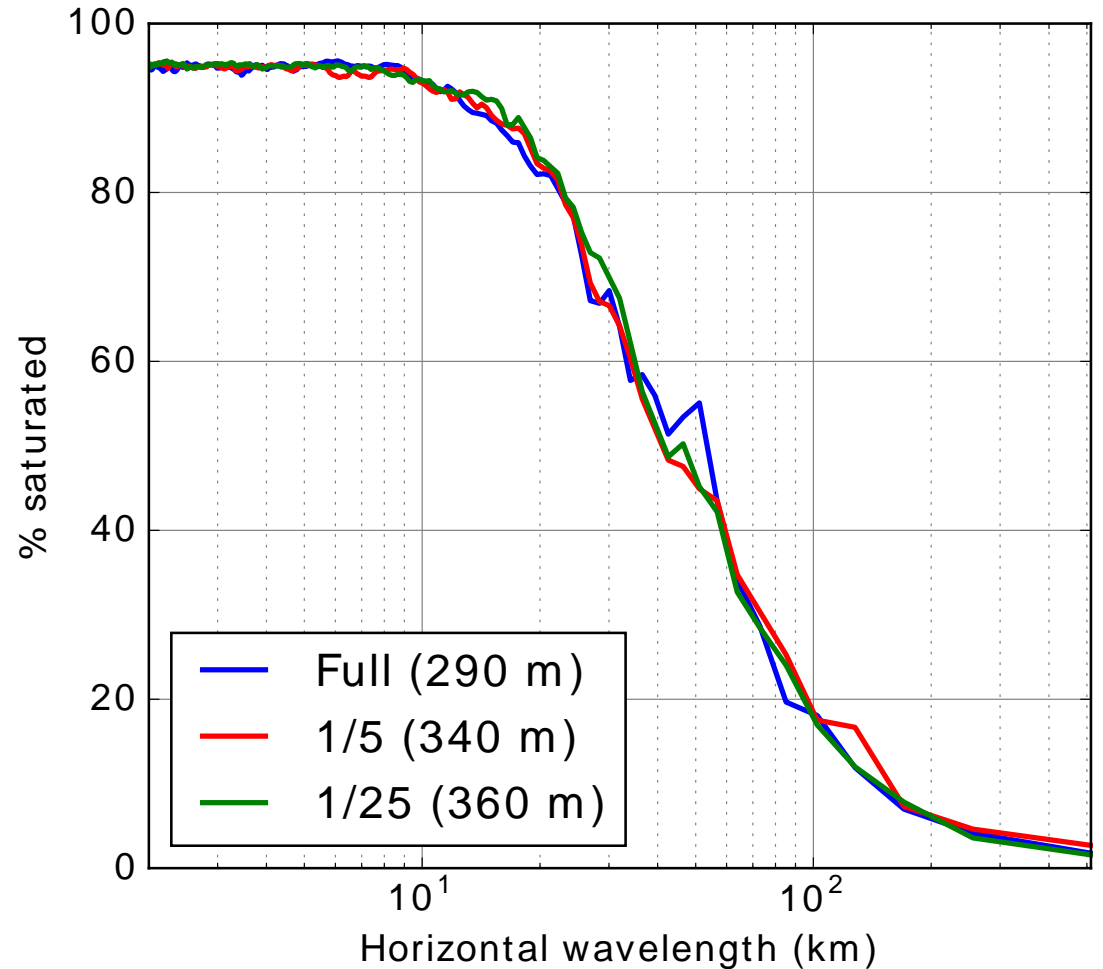
## What happens if the initial errors are lower amplitude?

- Multiply the magnitude of the  $0.1 \text{ g kg}^{-1}$  moisture perturbations by factors of  $1/5$  and  $1/25$
- Repeat the 8 and 128-km ensemble simulations for the  $20 \text{ m s}^{-1}$  shear case

# Similar predictability lead times from different initial amplitudes

Combined results for 8- and 128-km ensembles (40 members)

Gain 50 min, then just 20 min with the second reduction.



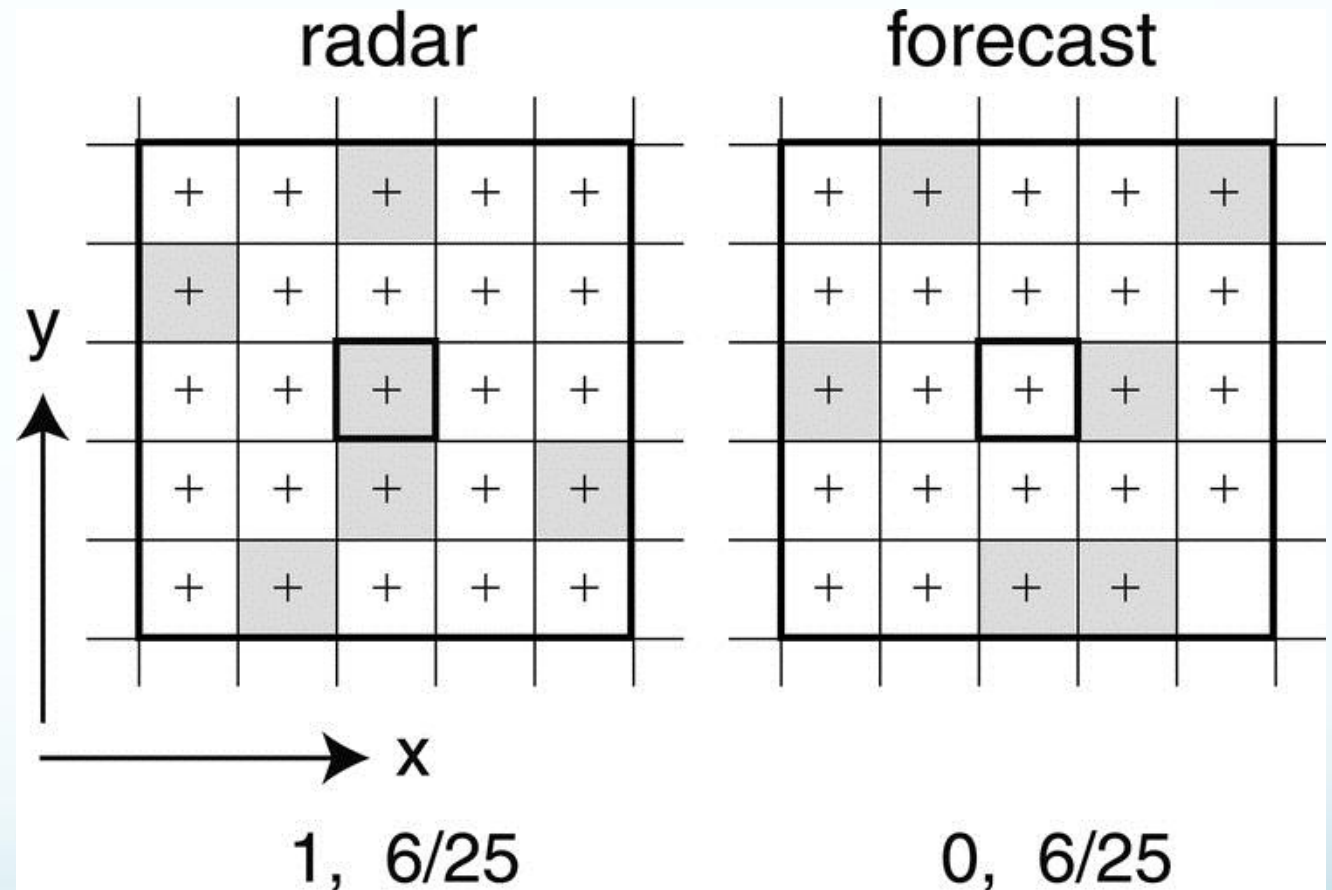
Would a forecaster view the three shear cases as having similar predictability?



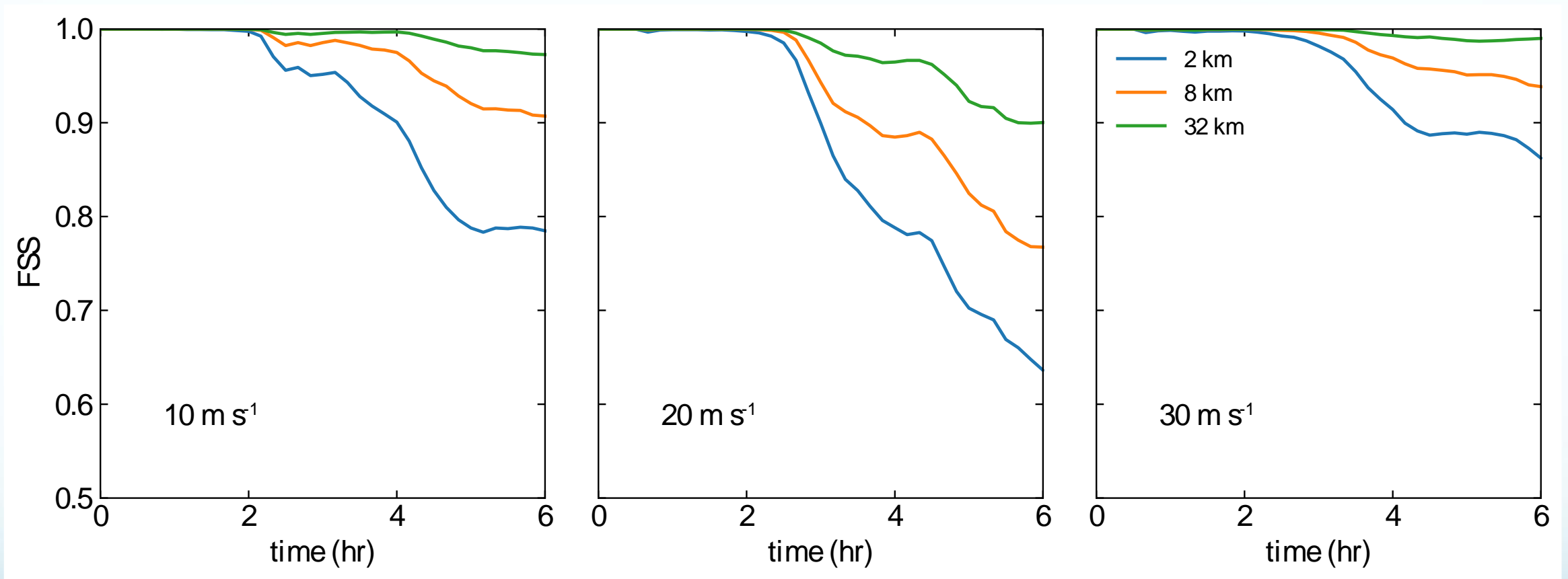
# Another measure of predictability

## Fractions skill score

(Roberts and Lean, *MWR*, 2008)



# FSS as a function of environmental shear

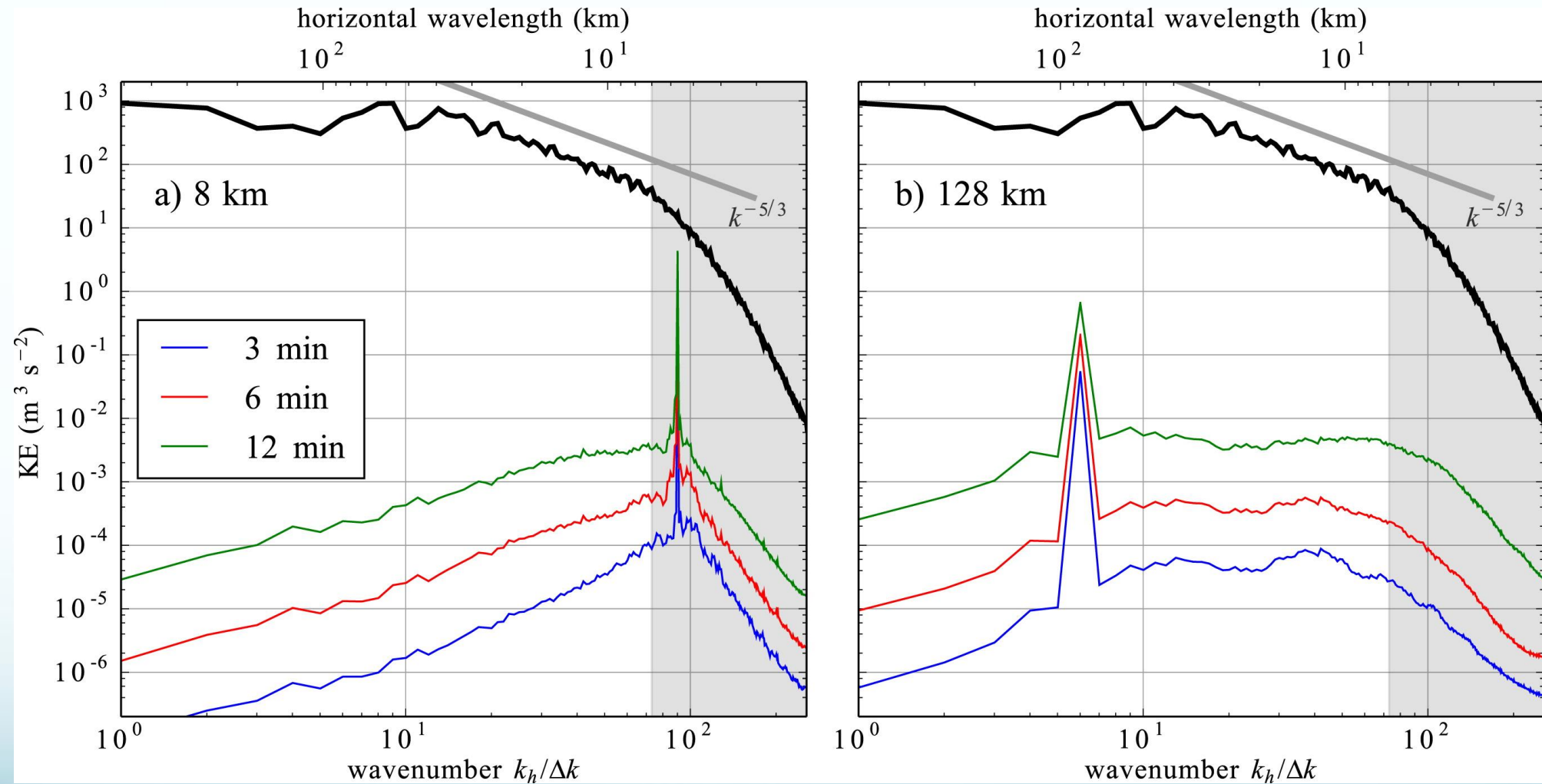


FSS score captures degree of variability suggested by visual inspection of the ensemble

# After large-scale errors spread downscale, is there a butterfly effect connected with subsequent upscale propagation?

- Add perturbations after 1 hour when the background KE spectrum has largely developed.
- 20 m s<sup>-1</sup> shear case
- Perturbation wavelengths
  - 8-km square wave has 2D wavelength 5.7 km
  - 128-km square wave has 2D wavelength 90 km

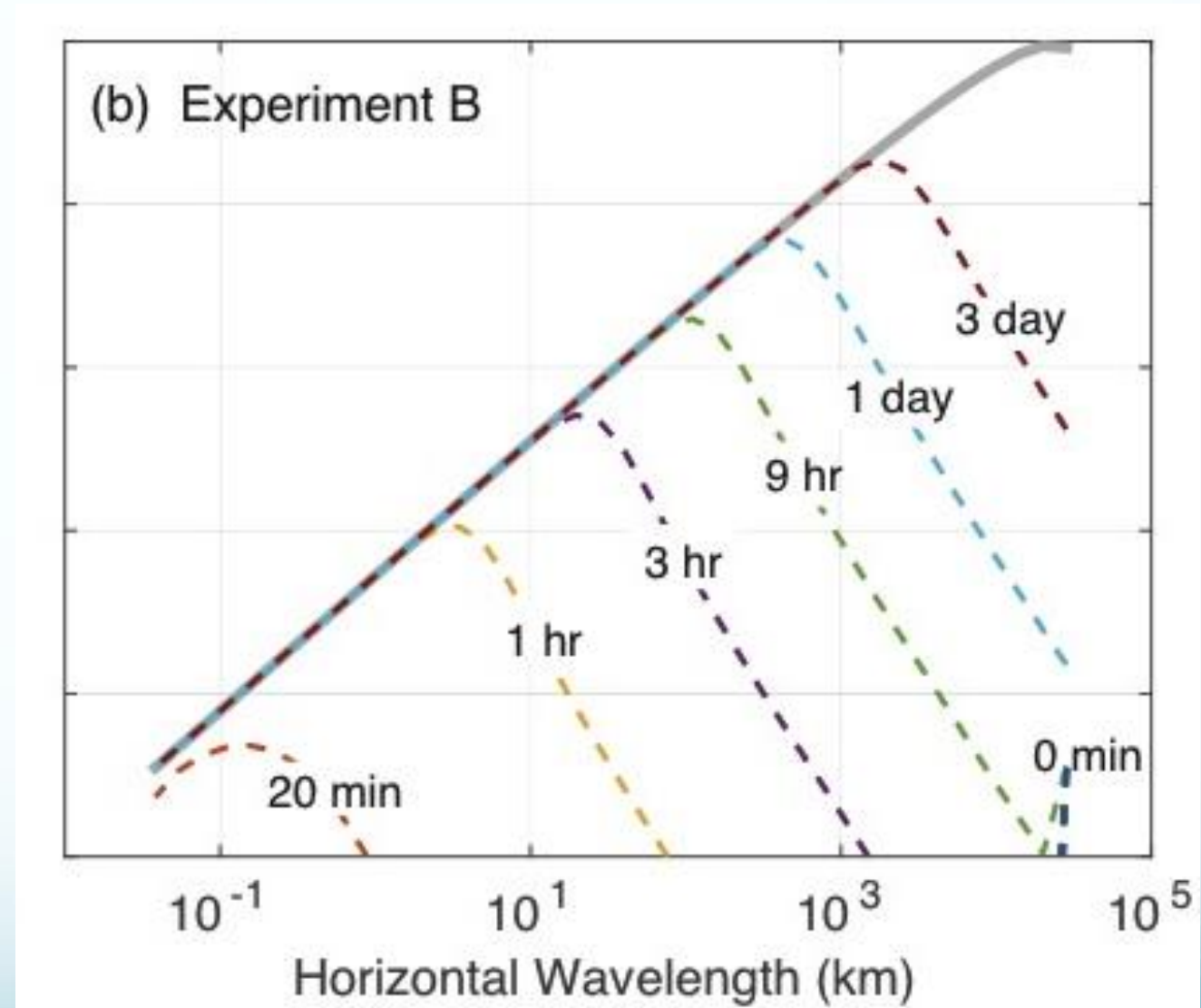
# Error growth at early times



Growth is primarily up amplitude.

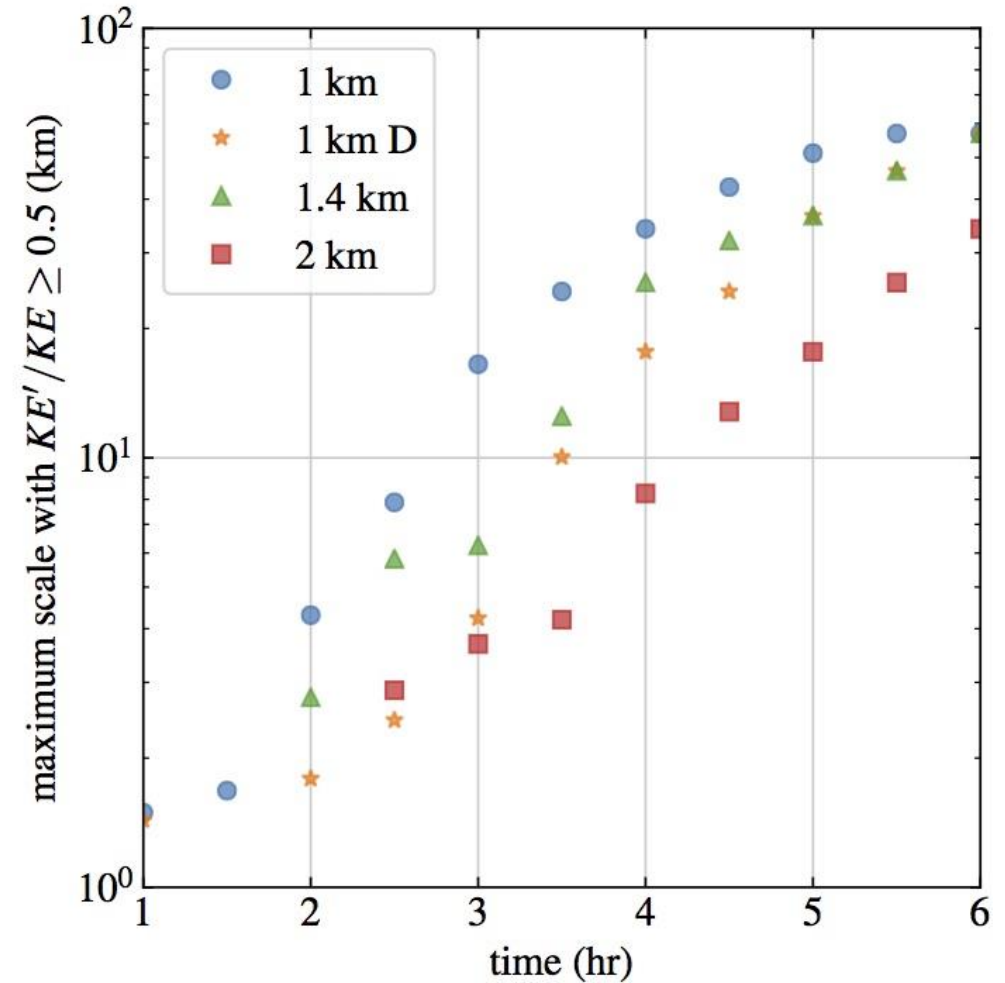
# Do the small scales still play an important role?

Upscale error growth after error propagates downscale in Lorenz's experiment B



# Upscale evolution of relative errors

- 25 m s<sup>-1</sup> environmental shear
- 1, 1.4 and 2 km grid spacing
- 1-km case with extra diffusion
- Initial errors are large scale (128 km square waves)
- Maximum wavelength for which  $KE'/KE > 0.5$  plotted as a function of time
- **Errors propagate more rapidly when small-scale motions are captured.**



# Implications for data assimilation: I

Parseval's relation

$$\int_S u^2(x) dx = \int_{-\infty}^{\infty} \hat{u}(k) \hat{u}^*(k) dk$$

KE in wavenumber band  $(k_1, k_2)$

$$E(k_1, k_2) = \int_{k_1}^{k_2} \hat{u}(k) \hat{u}^*(k) + \hat{v}(k) \hat{v}^*(k) dk$$

# Implications for data assimilation: II

- $k^{-5/3}$  KE spectrum

$$\frac{E(k_1, k_2)}{E(k_3, k_4)} = \frac{\lambda_1^{2/3} - \lambda_2^{2/3}}{\lambda_3^{2/3} - \lambda_4^{2/3}}$$

- Ratio of velocities in 200-400-km band to those in 2-4-km band is 0.21
- Which is the easier goal?
  - Reduce errors at 200-400 km below 10%
  - Reduce errors at 2-4 km below 50%



# Conclusions

- The large scales exert significant control on small-scale weather (Anthes), but **that control also includes the introduction of the many serious initial errors.**
- Small **relative** errors in the largest scales at which the background KE spectrum follows a  $k^{-5/3}$  slope (100--400 km) rapidly propagate down to the smallest resolved scale.
- Those small-scale errors subsequently propagate back upscale as if they had simply originated in the small scales.
  - Upscale growth is responsible for the finite limit to intrinsic predictability
- No easy way to diagnose the scale of the “original errors”.

# References

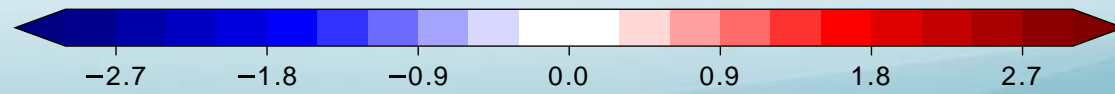
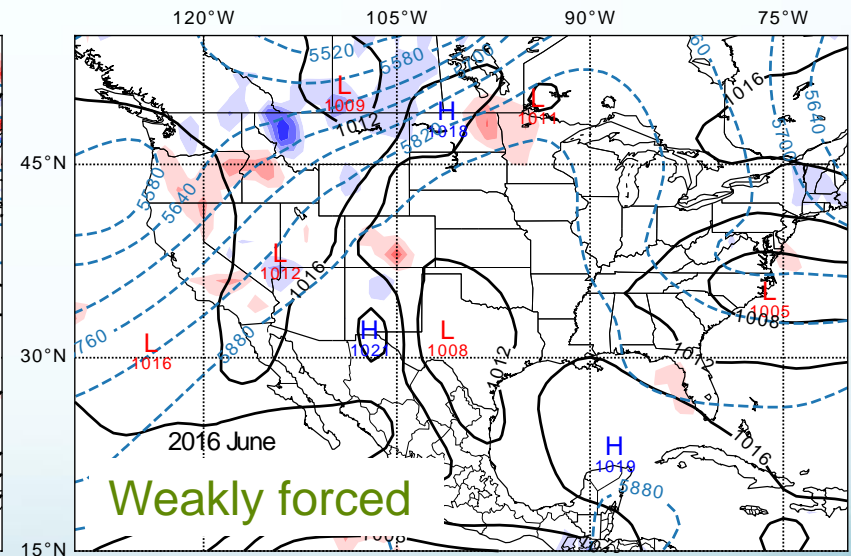
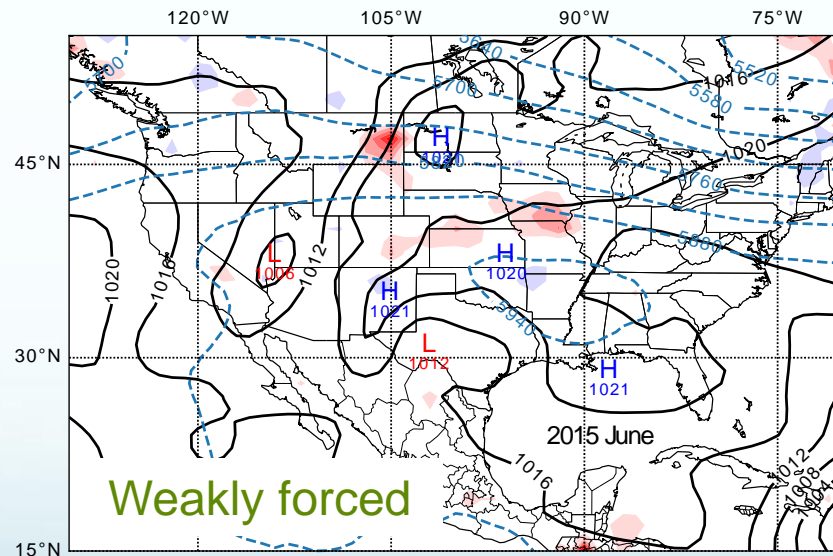
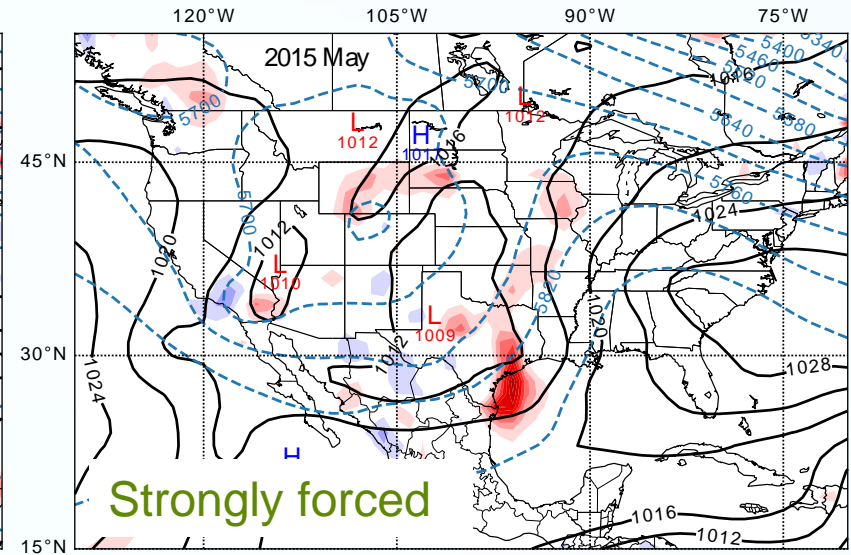
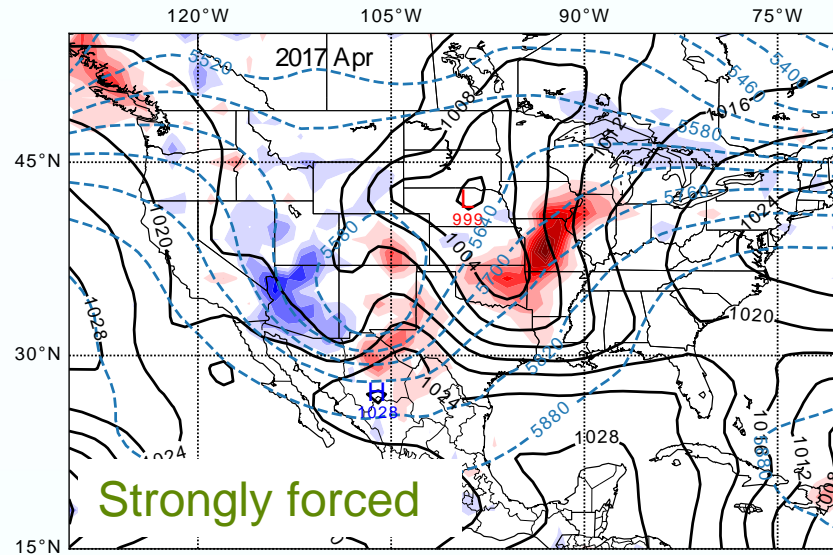
Weyn, J.A.. and D.R. Durran, 2017: The dependence of the predictability of mesoscale convective systems on the horizontal scale and amplitude of initial errors in idealized simulations. *J. Atmos. Sci.*, **74**, 2191-2210.

Durran, D.R. and J.A. Weyn, 2016: Thunderstorms don't get butterflies. *Bull. Amer. Meteor. Soc.*, **97**, 237-243.

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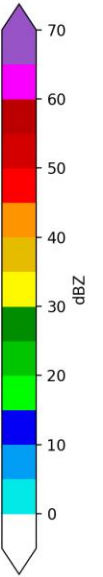
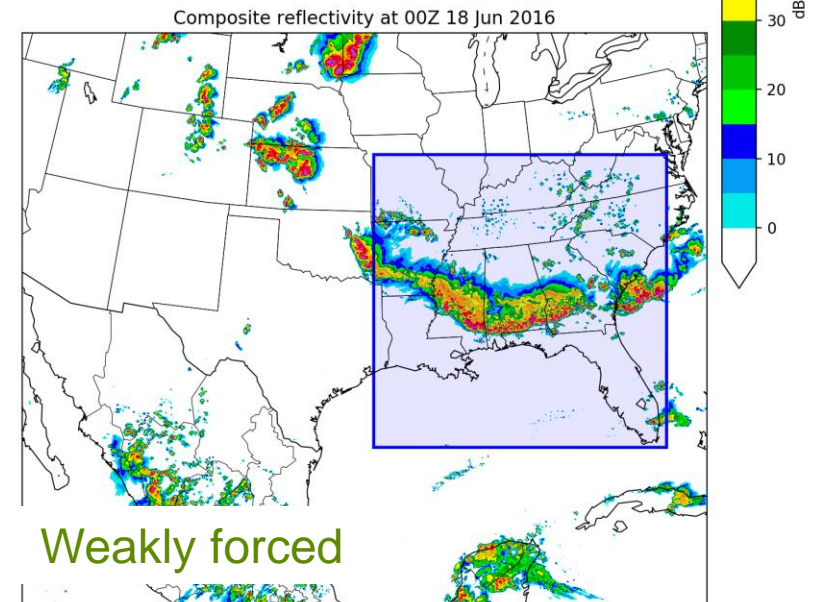
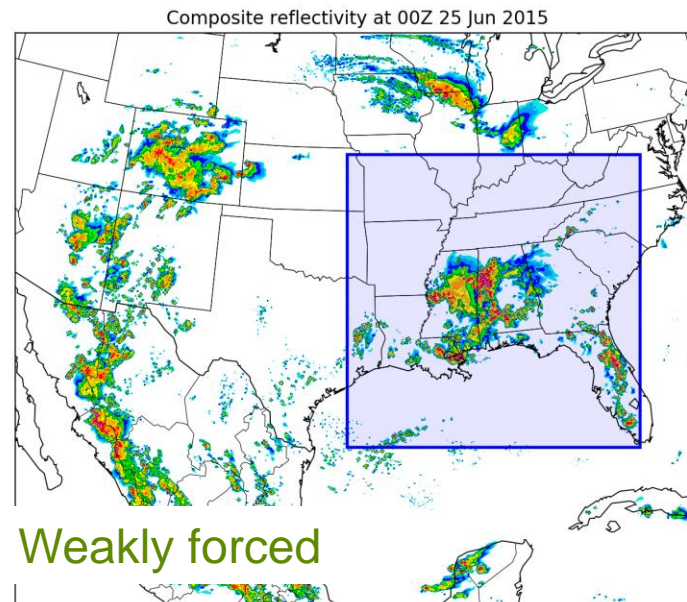
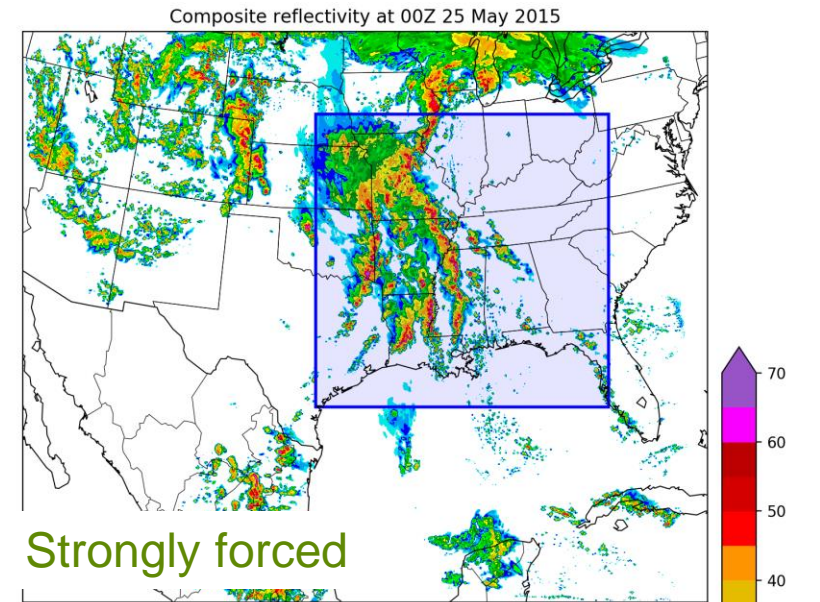
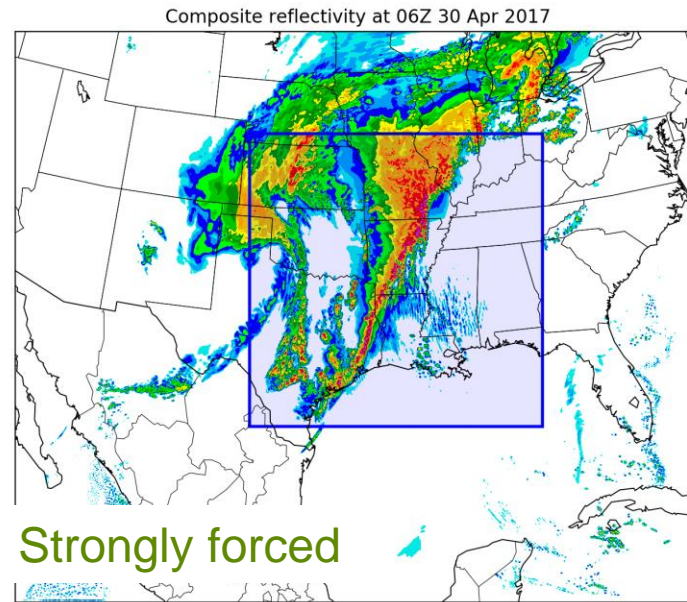
# Case studies

- Sea-level pressure
- 500 hPa heights
- 500 hPa vertical velocity (contours)



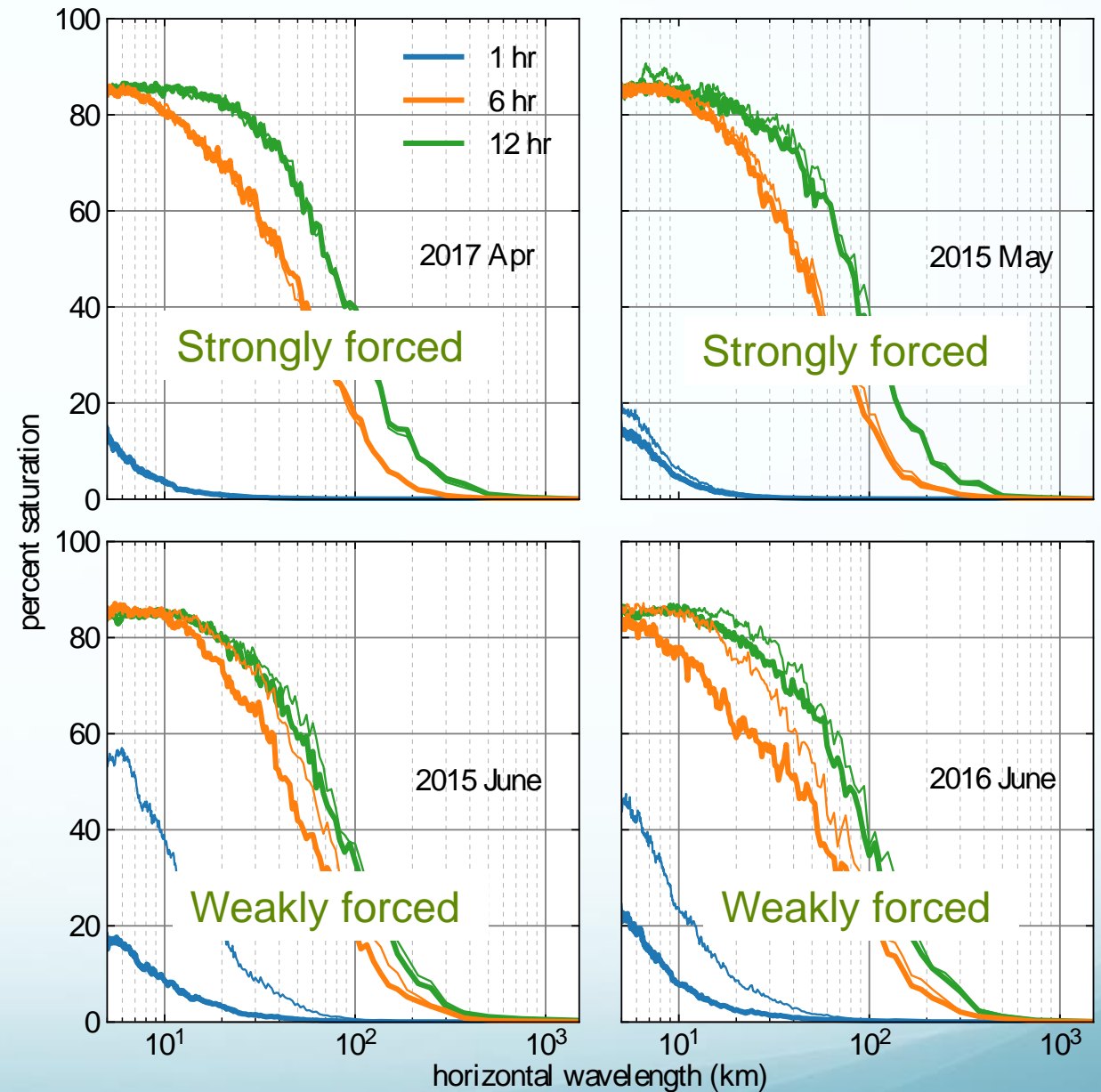
# Control Simulations

- Simulated composite reflectivity
- 6 other ensemble members
- 12 hours after initialization from GFS
- 20 or 200-km square wave moisture perturbations are introduced at hour 6
- 2.5 km horiz. resolution



# Error saturation ( $KE'/KE$ ) in layer $10 < z < 12$ km

- Similar errors at 12 hr in all cases
- Small-scale errors produce more saturation at 6 hr in the weakly forced cases
  - More variation in CI



# FSS: Case Studies

As in the  $KE'/KE$  saturation plots, in weakly forced cases:

20-km perturbations lead to larger early-time errors than 200-km perturbations.

